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## Fermions Destabilize Electroweak Strings

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*Z* strings in the Weinberg-Salam model including fermions are unstable for all values of the parameters. The cause of this instability is the fermion vacuum energy in the *Z*-string background. *Z* strings with nonzero fermion densities, however, may still be stable.

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The recent discovery of cosmic string solutions in the Weinberg-Salam model of electroweak interactions [1,2] has fueled a burst of activity in the study of defects in the standard model and its extensions, and of their possible consequences for astrophysics and cosmology [3,4]. The existence of these “electroweak strings,” which are essentially Nielsen-Olesen vortices embedded in the Weinberg-Salam model, was previously neglected because they do not possess topological stability. They may nonetheless be stable if they sit at a local minimum of the energy. Because they owe their existence to energetic rather than topological reasons, the stability of electroweak strings is sensitively dependent on the field content and the values of the parameters in the theory. For example, in a simplified version of the Weinberg-Salam model containing only bosonic fields, *Z* strings [2] are stable only for light Higgs masses ( $\leq m_Z$ ), and for  $\sin^2\theta_W$  fairly close to unity [5], a region that obviously does not include the physical world.

Attempts have been made to increase the range of stability of the *Z* string by extending the model to include other fields [3,6,7]. One idea, familiar from the study of nontopological solitons, is to include particles whose mass arises from the Higgs mechanism. Such particles remain massless at the center of the string where the Higgs field vanishes, and the presence of such particles at the core would resist the string’s dissolution, because that would increase their energy. Indeed, the presence of charged scalar bound states was shown to lower the value of  $\sin^2\theta_W$  for which the string is stable [6].

It has been suggested that a similar enhancement of stability could be achieved by using fermion bound states on the *Z* string [6]. The existence of *Z*-string zero modes, fermion states localized on the string with zero energy, lends support to this idea [8–10]. Another advantage of this suggestion is that fermions are already contained in the standard electroweak model.

In this Letter, we show that, on the contrary, the presence of fermions in the electroweak theory *destabilizes* *Z* strings. More precisely, the lowest-energy (or ground) state of the *Z* string is always a local maximum of the energy functional with respect to (at least) one of the modes of instability. The *Z* string is therefore unstable for *all* values of the parameters of the Weinberg-Salam model. (It

is possible, however, that a higher-energy state of the *Z* string, with a finite quark density, could be locally stable.)

This instability results from the fermion vacuum energy, which also has an important effect on other types of solitons [11]. One cannot consistently consider the effects of positive-energy fermion states without also taking account of the (filled) negative-energy states, particularly because, with the existence of zero modes, there is no gap between them. We will show that the contribution to the energy functional of the filled Dirac sea, i.e., the fermion vacuum energy, is a local maximum for the *Z* string.

First, we will describe the fermion spectrum in the presence of the *Z* string; then we will show how the fermion vacuum energy changes under certain small perturbations away from the *Z* string. The electroweak Lagrangian is

$$L = L_{\text{boson}} + \sum L_{\text{quark}} + \sum L_{\text{lepton}}, \quad (1)$$

with

$$L_{\text{boson}} = -\frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + |D_\mu^L \Phi|^2 - \lambda \left( \Phi^\dagger \Phi - \frac{\eta^2}{2} \right)^2, \quad (2)$$

where  $W_{\mu\nu}^a$  and  $F_{\mu\nu}$  are field strength tensors for the  $SU(2)_L$  and  $U(1)_Y$  gauge fields  $W_\mu^a$  and  $B_\mu$ , respectively, and  $\Phi = \begin{pmatrix} \phi_1 \\ \phi_0 \end{pmatrix}$  is the complex Higgs doublet. Each quark doublet contributes a term

$$L_{\text{quark}} = \bar{\Psi}^L i \not{D}^L \Psi^L + \bar{\psi}_+^R \not{D}^R \psi_+^R + \bar{\psi}_-^R i \not{D}^R \psi_-^R - G_+ (\bar{\Psi}^L \tilde{\Phi} \psi_+^R + \bar{\psi}_+^R \tilde{\Phi}^\dagger \Psi^L) - G_- (\bar{\Psi}^L \Phi \psi_-^R + \bar{\psi}_-^R \Phi^\dagger \Psi^L), \quad (3)$$

where  $\Psi^L = \begin{pmatrix} \psi_+^L \\ \psi_-^L \end{pmatrix}$  and  $\tilde{\Phi} = i\tau^2 \Phi^* = \begin{pmatrix} \phi_0^* \\ -\phi_1^* \end{pmatrix}$ . Each lepton doublet contributes the same term, absent any pieces containing  $\psi_+^R$ . (We neglect interfamily mixing.) The gauge-covariant derivatives in Eqs. (1) and (2) are

$$D_\mu^L = \partial_\mu - \frac{ig}{2} \tau^a W_\mu^a - \frac{ig'}{2} Y B_\mu, \\ D_\mu^R = \partial_\mu - \frac{ig'}{2} Y B_\mu, \quad (4)$$

with  $Y$  the hypercharge of the field on which the covariant derivative is acting.

The  $Z$  string [2] is the field configuration

$$\Phi = \frac{\eta f(\rho)}{\sqrt{2}} e^{i\phi} \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \quad \begin{pmatrix} Z^1 \\ Z^2 \end{pmatrix} = \frac{2v(\rho)}{\alpha\rho} \begin{pmatrix} -\sin\phi \\ \cos\phi \end{pmatrix}, \quad (5)$$

all other fields vanishing, where  $f(\rho)$  and  $v(\rho)$  obey the Nielsen-Olesen equations [12]

$$f'' + \frac{f'}{\rho} - (1-v)^2 \frac{f}{\rho^2} + \lambda\eta^2(1-f^2)f = 0, \\ v'' - \frac{v'}{\rho} + \frac{\alpha^2\eta^2}{4} f^2(1-v) = 0, \quad (6)$$

with boundary conditions

$$f(0) = v(0) = 0, \quad f(\rho) \xrightarrow{\rho \rightarrow \infty} 1, \quad v(\rho) \xrightarrow{\rho \rightarrow \infty} 1. \quad (7)$$

Recall that  $Z_\mu = \cos\theta_W W_\mu^3 - \sin\theta_W B_\mu$  and  $\alpha = \sqrt{g^2 + g'^2}$ .

The main question is whether such a field configuration is energetically stable. One possible mode of instability is that the upper component  $\phi_1$  of the Higgs field may develop a nonzero value, allowing the  $Z$  string to unwind. (There may be other modes of instability as well.) To determine whether the string is stable to small perturbations in this direction (in the absence of fermions), one computes the change in the bosonic field energy

$$\Delta E_{\text{boson}}[\phi_1] = E_{\text{boson}}[f, v; \phi_1] - E_{\text{boson}}[f, v; 0]. \quad (8)$$

If the  $Z$  string is stable,  $E_{\text{boson}}[f, v; 0]$  is a local minimum of the energy functional, and  $\Delta E_{\text{boson}}[\phi_1]$  will be quadratic in  $\phi_1$  with a positive coefficient [2].

Next, we consider the effect of fermions in the electroweak theory. The Dirac equation in the  $Z$  string background has the form

$$\gamma^\mu \left( i\partial_\mu - \frac{\alpha\ell_\pm}{2} Z_\mu \right) \psi_\pm^L - m_\pm f(\rho) e^{\mp i\phi} \psi_\pm^R = 0, \\ \gamma^\mu \left( i\partial_\mu - \frac{\alpha r_\pm}{2} Z_\mu \right) \psi_\pm^R - m_\pm f(\rho) e^{\pm i\phi} \psi_\pm^L = 0, \quad (9)$$

where  $m_\pm = G_\pm \eta / \sqrt{2}$ ,  $\ell_\pm = (y \pm 1) \sin^2\theta_W \mp 1$ , and  $r_\pm = (y \pm 1) \sin^2\theta_W$ , with  $y$  the hypercharge of the left-handed doublet  $\Psi^L$ . This equation has zero-energy modes [8-10,13], which obey  $\gamma_0\gamma_3\psi_\pm = \pm\psi_\pm$ . Using the chiral representation for the Dirac matrices

$$\gamma^5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \gamma^0 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}, \\ \gamma^i = \begin{pmatrix} 0 & \tau^i \\ -\tau^i & 0 \end{pmatrix}, \quad (10)$$

and recalling that  $\psi_\pm^L = \frac{1}{2}(1 - \gamma^5)\psi_\pm$  and  $\psi_\pm^R = \frac{1}{2}(1 + \gamma^5)\psi_\pm$ , one may write the zero mode solutions as

$$\psi_{\pm,0} = e^{\pm \int_0^{\rho} [\ell_\pm v(\rho')/\rho'] d\rho'} \begin{pmatrix} \frac{i}{m_\pm f(\rho)} P'_\pm(\rho) \chi_\pm \\ P_\pm(\rho) \chi_\mp \end{pmatrix}, \quad (11)$$

where  $\chi_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ,  $\chi_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ , with  $P_\pm(\rho)$  obeying the equation

$$P''_\pm - \frac{f'}{f} P'_\pm \pm \frac{(\ell_\pm + r_\pm)v}{\rho} P'_\pm - m_\pm^2 f^2 P_\pm = 0 \quad (12)$$

and normalized by

$$\int d^3x e^{\pm 2 \int_0^\rho [\ell_\pm v(\rho')/\rho'] d\rho'} [P_\pm^2 + (P'_\pm/m_\pm f)^2] = 1. \quad (13)$$

For the neutrino ( $m_+ = 0, \ell_+ = -1$ ), Eq. (12) has the simple solution  $P_+ = 1$ , but by Eq. (7), for  $\rho \rightarrow \infty$ ,  $\psi_{+,0} \rightarrow 1/\rho$ , so the zero mode is not normalizable, at least for a straight infinite string (but see Ref. [9]). Equation (12) has the explicit solution

$$P_\pm(\rho) = N e^{-\int_0^\rho m_\pm f(\rho') d\rho'} \quad (14)$$

for the special case  $y = 0$  and  $\cos^2\theta_W = \frac{1}{2}$ .

The existence of zero modes generates a  $2^N$ -fold degeneracy of the  $Z$ -string ground state, where  $N$  is the number of quark and charged lepton flavors. The ground state of the string will have the global quantum numbers of each fermion flavor [14], either  $\frac{1}{2}$  or  $-\frac{1}{2}$ , depending on whether the corresponding zero mode is occupied or not. The occupation of the zero modes will not alter the Nielsen-Olesen equations (6) for the string profile, because the fermion source term for the  $\phi_0$  and  $Z_\phi$  fields vanishes for the zero modes.

In the  $(3+1)$ -dimensional context of the  $Z$  string, the zero-energy solution (11) generates a whole family of solutions of the Dirac equation

$$\psi_{\pm,p}(\rho, z, t) = e^{ipz - i\epsilon_{\pm,p}t} \psi_{\pm,0}(\rho) \quad (15)$$

with energies

$$\epsilon_{\pm,p} = \pm p. \quad (16)$$

These solutions correspond to massless chiral fermions confined to the  $Z$  string; the up-type quarks run up the string (in the  $+z$  direction) and the down-type quarks and charged leptons run down the string (in the  $-z$  direction) at the speed of light. In addition to these "massless" solutions of the Dirac equation, there are many "massive" solutions, whose energies are separated from zero by a finite gap.

What effect do fermions have on the stability of the  $Z$  string? Earnshaw and Perkins [8] pointed out that the fermion zero mode provides a nonvanishing source term in the equation of motion for  $\phi_1$ . This violates the "Vachaspati existence criterion" [2] and would appear to imply that the  $Z$ -string configuration with  $\phi_1 = 0$  is not an extremum of the energy. Such a conclusion, however, would be premature.

The reason that the zero modes are a source for  $\phi_1$  is that the presence of a nonzero value of  $\phi_1$  lifts the degeneracy between the  $\psi_{+,0}$  and  $\psi_{-,0}$  zero modes, one linear combination of the zero modes shifting up and the orthogonal combination shifting down. The lower eigenstate is

filled in the ground state of the  $Z$  string, so its descent lowers the  $Z$ -string energy. Before drawing any conclusions about the overall stability of the  $Z$  string, however, we must determine the effect of the  $\phi_1$  perturbation on the rest of the fermion eigenenergies.

The effective energy of the  $Z$ -string ground state is

$$E_{\text{eff}} = E_{\text{boson}} + E_{\text{fermion}}, \quad (17)$$

where  $E_{\text{boson}}$  is the bosonic field energy and  $E_{\text{fermion}}$  the fermion vacuum energy in the  $Z$ -string background (i.e., the energy of the filled Dirac sea). The change  $\Delta E_{\text{boson}}[\phi_1]$  under a small perturbation  $\phi_1$  was considered above (8); the change in the fermion vacuum energy

is

$$\Delta E_{\text{fermion}}[\phi_1] = \sum_{\epsilon_{+,n} < 0} \Delta \epsilon_{+,n}[\phi_1] + \sum_{\epsilon_{-,n} < 0} \Delta \epsilon_{-,n}[\phi_1] + \delta E[\phi_1], \quad (18)$$

where  $\Delta \epsilon_{\pm,n}[\phi_1]$  denotes the shift of the  $Z$ -string Dirac eigenenergies  $\epsilon_{\pm,n}$  under the perturbation, and the sum is over negative-energy eigenvalues only.

We compute the eigenvalue shifts  $\Delta \epsilon_{\pm,n}[\phi_1]$  perturbatively in  $\phi_1$ . Because the perturbation is off-diagonal in the  $+$  and  $-$  fields, the leading shift is second order, and the change in fermion vacuum energy is

$$\begin{aligned} \Delta E_{\text{fermion}}[\phi_1] = & \sum_{\epsilon_{+,n} < 0} \sum_{\epsilon_{-,m} > 0} \frac{|\int d^3x (G_- \bar{\psi}_{-,m}^R \psi_{+,n}^L - G_+ \bar{\psi}_{-,m}^L \psi_{+,n}^R) \phi_1^*|^2}{\epsilon_{+,n} - \epsilon_{-,m}} \\ & + \sum_{\epsilon_{-,n} < 0} \sum_{\epsilon_{+,m} > 0} \frac{|\int d^3x (G_- \bar{\psi}_{+,m}^L \psi_{-,n}^R - G_+ \bar{\psi}_{+,m}^R \psi_{-,n}^L) \phi_1|^2}{\epsilon_{-,n} - \epsilon_{+,m}} + \delta E[\phi_1]. \end{aligned} \quad (19)$$

The sums over intermediate energy eigenstates  $\epsilon_{\pm,m}$  include only positive-energy states; the contributions from negative-energy intermediate states cancel between the two sums. The sums in Eq. (19) diverge in the ultraviolet. The  $Z$  string is not responsible for this, because the same divergence occurs in the usual constant field background. In that case, the divergence is canceled by adding a counterterm  $\delta E[\phi_1]$ . The same counterterm will suffice to render  $\Delta E_{\text{fermion}}[\phi_1]$  ultraviolet finite.

Let us evaluate the shifts in eigenenergies corresponding to the massless solutions (15) more explicitly. First, restrict the perturbation  $\phi_1$  to a constant (complex) value  $\eta g/\sqrt{2}$  over the region where the zero mode wave function (11) is appreciable (but let  $\phi_1 \rightarrow 0$  as  $\rho \rightarrow \infty$ ). Second, noting that  $E_{\text{fermion}}$  is proportional to the length of the string (as is  $E_{\text{boson}}$ ), consider a  $Z$  string of length  $L$ . Periodic boundary conditions on the fermion wave functions restrict the  $z$  momenta to  $p = 2\pi n/L$ . Taking  $L$  large, the sum of the energy shifts of the massless states becomes

$$\Delta E'_{\text{fermion}}[g] = -\frac{L}{\pi} |g|^2 |A|^2 \int_0^\Lambda \frac{dp}{2p}, \quad (20)$$

with

$$\begin{aligned} A = & 2\pi iL \int \rho d\rho \frac{(P_+ P_-)'}{f} \\ & \times \exp\left(\int_0^\rho [(\ell_+ - \ell_-)v(\rho')/\rho'] d\rho'\right), \end{aligned} \quad (21)$$

where we have included only massless intermediate energy states. Among this subset of intermediate states, a selection rule ensures that the perturbation only couples the eigenstates  $\psi_{\pm,p}$  to the eigenstate  $\psi_{\mp,-p}$ .

The integral over momenta (20) diverges both in the ultraviolet and in the infrared. As previously noted, the ultraviolet divergence is canceled by counterterms;

the infrared divergence signals the breakdown of the perturbative expansion when the energy denominator  $2p$  becomes smaller than the perturbation  $g$ . We redo the calculation for states with small  $p$ , now treating  $p$  as part of the perturbation. The unperturbed states are now degenerate; degenerate perturbation theory yields the perturbed energies

$$\begin{vmatrix} \epsilon - p & gA \\ g^* A^* & \epsilon + p \end{vmatrix} = 0 \implies \epsilon = \pm \sqrt{p^2 + |gA|^2}. \quad (22)$$

As mentioned above, the degenerate zero mode ( $p = 0$ ) states are resolved into  $\epsilon = \pm |gA|$ . This improved calculation yields an infrared-finite result

$$\begin{aligned} \Delta E'_{\text{fermion}}[g] = & -\frac{L}{2\pi} \int_{-\Lambda}^\Lambda dp \left( \sqrt{p^2 + |gA|^2} - |p| \right) \\ \xrightarrow{\Lambda \rightarrow \infty} & -\frac{L}{4\pi} |gA|^2 \left[ 1 + \ln\left(\frac{4\Lambda^2}{|gA|^2}\right) \right]. \end{aligned} \quad (23)$$

The ultraviolet divergence is absorbed by the counterterm, leading to a completely finite expression for the change in the fermion vacuum energy (per unit length) under the perturbation  $\phi_1$ :

$$\Delta E_{\text{fermion}}[g] = L \left( \frac{|A|^2}{4\pi} |g|^2 \ln|g|^2 + C_f |g|^2 \right). \quad (24)$$

The coefficient  $C_f$  receives contributions from the shifts of the massive Dirac eigenvalues as well as from the finite part of the counterterm. Each quark doublet contributes a term of the form (24) to the fermion vacuum energy (with different values of  $A$ ); the charged leptons do not contribute because their eigenenergies are not shifted by the perturbation (in the absence of normalizable neutrino zero modes).

The change in the bosonic field energy for the perturbation we are considering has the form [2]

$$\Delta E_{\text{boson}}[g] = LC_b |g|^2, \quad (25)$$

where the sign of  $C_b$ , which depends on the parameters of the model, determines whether the bosonic  $Z$  string (without fermions) is stable in the direction of this perturbation. Thus, the effective energy of the ground state of the  $Z$  string is

$$E_{\text{eff}}[g] = E_{\text{eff}}[0] + L \left( C |g|^2 + \frac{1}{4\pi} |g|^2 \ln |g|^2 \sum_{\text{quark}} |A|^2 \right). \quad (26)$$

Observe first that  $|g| = 0$  is an extremum of this expression, so even in the presence of fermions the  $Z$ -string configuration (5) with  $|\phi_1| = 0$  remains a solution of the equations of motion. This extremum, however, is necessarily a maximum, regardless of the value of  $C$ . (See Ref. [15] for a similar phenomenon in two dimensions.) Hence, the  $Z$  string is unstable to perturbations in  $\phi_1$  for *all* values of the parameters of the Weinberg-Salam model.

For a  $Z$  string of large but finite length  $L$ , the fermion energy (23) becomes a sum over momenta, given by

$$\Delta E'_{\text{fermion}}[g] = -|gA| - \frac{L|gA|^2}{2\pi} \left[ \ln \left( \frac{L\Lambda}{2\pi} \right) + \gamma \right] + \dots$$

in the limit  $g \ll 1/L$ . The leading term linear in  $g$  can be canceled by populating both (or neither) zero modes. The subleading term then contributes energy per unit length quadratic in  $g$  with a negative coefficient that diverges as  $\ln L$ . Thus, regardless of the bosonic contribution (25), the  $Z$ -string ground state is unstable for large  $L$ , the same conclusion reached above.

This does not mean that no stable nontopological string configuration exists. A (locally) stable string with  $\phi_1$  slightly displaced from zero may exist, though that remains to be demonstrated. What we are saying is that the simple Nielsen-Olesen string embedded into the Weinberg-Salam model with all other fields vanishing is necessarily unstable.

Most attempts to increase the stability of the  $Z$  string do so by increasing the coefficient  $C$  in (26); this will not work here since no coefficient, however positive, can outweigh the negative curvature at  $|g| = 0$  caused by the  $|g|^2 \ln |g|$  term. The only way to overcome this term is to occupy some of the *positive*-energy fermion states. This requires not only a single fermion (as in the case of nontopological solitons) but rather a finite density of

fermions along the  $Z$  string. If the string holds  $\zeta$  positive-energy fermions per unit length of types  $+$  and  $-$ , the effective energy will change by

$$\Delta E[g] = \frac{L}{2\pi} \int_{-2\pi\zeta}^{2\pi\zeta} dp \left( \sqrt{p^2 + |gA|^2} - |p| \right) \xrightarrow{|gA| \ll 2\pi\zeta} \frac{L}{4\pi} |gA|^2 \left[ 1 + \ln \left( \frac{16\pi^2 \zeta^2}{|gA|^2} \right) \right]. \quad (27)$$

If the  $Z$  string carries a nonzero density of quarks of *each* flavor, the change in energy will cancel the  $|g|^2 \ln |g|^2$  piece in (26), rendering the total energy proportional to  $|g|^2$ . A difficult calculation would then be required to determine the minimum fermion density necessary to stabilize the state. Even so, the state might only be a local minimum, liable to decay to a nonstring state via tunneling or thermal fluctuations.

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